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SUMMARY

The XV-15 Tilt-Rotor wing has six major aeroelastic modes that are close in frequency. To precisely excite individual modes during flight test, dual flaperon exciters with automatic frequency-sweep controls were installed. The resulting structural data were analyzed in the frequency domain (Fourier-transformed) with cross-spectral and transfer-function methods. Modal frequencies and damping were determined by performing curve fits to transfer-function magnitude and phase data and to cross-spectral magnitude data. Results are given for the XV-15 with its original metal rotor blades. Frequency and damping values are also compared with earlier predictions.

INTRODUCTION

Distinctive features of the XV-15 Tilt Rotor are the large wing-tip pylons which house the engines, transmissions and pivoting mechanisms for each rotor (fig. 1). The highest speeds are obtained in the cruise mode, making it the critical operating mode for aeroelastic stability. The concentrated masses at the wing tips keep the modal frequencies fairly low. Also, aeroelastic coupling between each rotor and pylon is destabilizing. Consequently, close attention must be paid to potential whirl-mode flutter during flight test. The problem is not unique to the XV-15 research aircraft, but is fundamental to any tilt-rotor aircraft of similar configuration, such as the XV-3, for which extensive studies were done (refs. 1 and 2), and the upcoming V-22 Osprey. The impact of aeroelastic stability requirements on tilt-rotor design is discussed in reference 3.

The planned flight tests of new composite rotors on the XV-15 (ref. 4) provided the incentive for a thorough reevaluation of XV-15 aeroelastics using the latest flight-test and modal identification techniques. The major wing modes were excited with flaperon frequency sweeps. Modal frequencies and damping were determined by performing curve fits to transfer-function magnitude and phase data and to cross-spectral magnitude data. The analysis programs have been used successfully on other flight data, notably for XV-15 aircraft flight dynamics (ref. 5).

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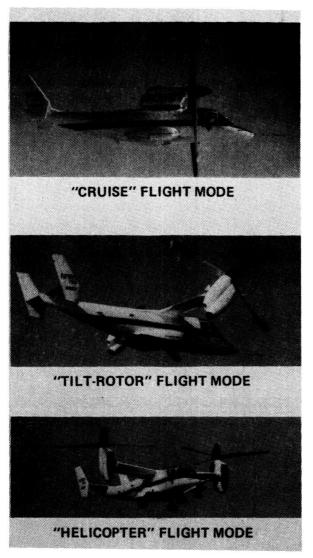


Figure 1.- XV-15 tilt-rotor research aircraft.

More complete discussions of the aeroelastic behavior of the XV-15, including previous flight-test data, are given in references 6 and 7. All flight-test data discussed in the present report are for the original metal blades and new steel hubs. References 6 and 7 give data for the metal blades and original titanium hubs.

This report presents discussions of the XV-15 aeroelastic modes and the flight-test techniques used to excite them; the analytical procedures used to extract modal frequencies and damping from flight-test data; and plots of estimated frequency and damping versus airspeed, including comparisons with theoretical values.

FLIGHT-TEST METHODS

The six major wing/pylon aeroelastic modes are illustrated in figures 2 and 3 and are listed in figure 4, which shows the ranges of natural frequencies identified during the flight-test program. Note that some predictions were inaccurate, even

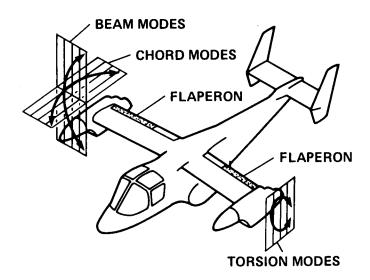


Figure 2.- XV-15 aircraft, showing flaperons and wing modes.

though they include refinements made after the aircraft was initially tested (ref. 6). The predictions and experimental data together indicate that at least one mode lies within the design rotor-speed range. Consequently, the rotor speed is restricted in cruise to 8.6 Hz (86% of 601 rpm) instead of the design minimum of 7.6 Hz. Several modes are very close; two--the antisymmetric-chord and antisymmetric-torsion modes--actually overlap, which was not predicted. It is this close placement in frequency and restricted rotor-speed range that makes precise identification of individual modes both difficult and necessary.

In earlier flight tests (refs. 6 and 7), frequency and damping were identified by using primarily the exponential-decay technique with a Prony analysis (ref. 8). The natural frequencies were precisely identified, but the damping values showed a great deal of scatter, especially where neighboring modes were grouped closely together. Limited data were obtained using turbulence excitation and analyzed with RANDOMDEC techniques (ref. 6), but again with considerable scatter. Frequency sweeps were also tried, but the original exciters were unable to adequately excite the modes (ref. 7).

Additional flight tests (ref. 5) used turbulence excitation, followed by frequency-domain analyses to identify frequencies and damping. The turbulence did not excite the modes strongly enough in most cases. However, the results of those tests indicated that the frequency-domain method was the most promising approach, provided that dual flaperon exciters were used to selectively excite the symmetric and antisymmetric modes, as recommended in reference 6. Compared to the exponential-decay and RANDOMDEC methods, the frequency-domain method is less sensitive to noise and makes more efficient use of flight time; accordingly, it is the method of choice.

The earlier flight tests used a high-frequency, limited-authority servo actuator in series with the right-flaperon control linkage to excite the bending and

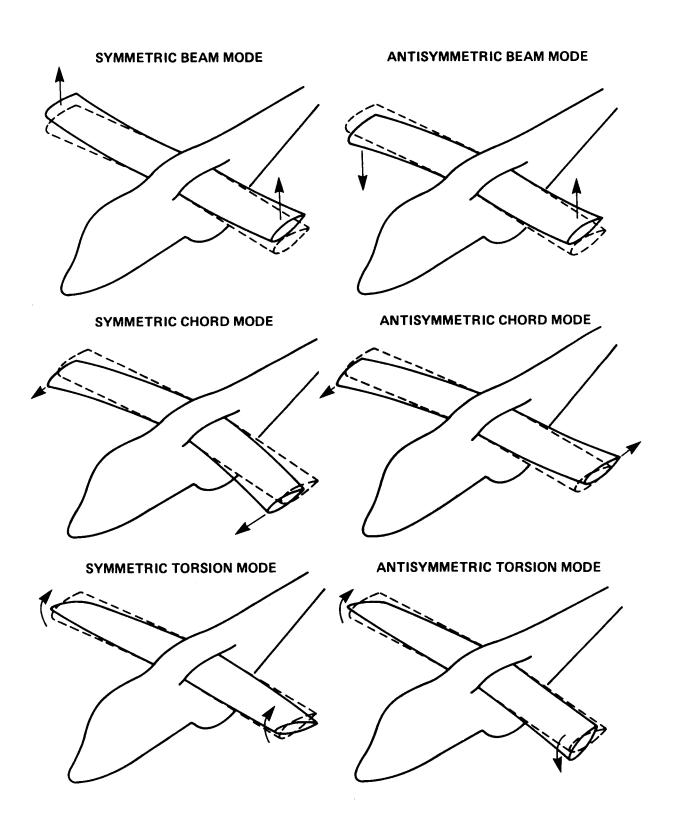


Figure 3.- XV-15 aeroelastic wing modes.

O PREDICTED FREQUENCY (NASTRAN) MEASURED FREQUENCY (NAVFIT)

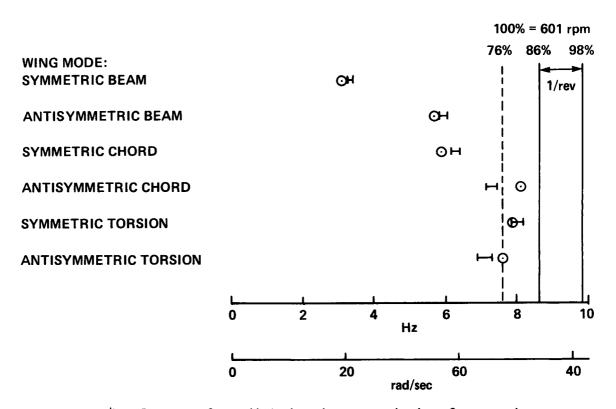


Figure 4.- Ranges of predicted and measured wing frequencies.

torsion modes, and a similar actuator in series with the right-rotor collective control to excite the chord modes. By adding a flaperon exciter to the left wing (fig. 2), different modes could be separately excited more easily. Symmetric modes were excited by driving the flaperons in phase; antisymmetric modes were excited by driving the flaperons in opposite phase. Furthermore, the paired flaperon exciters adequately excited the chord modes without using the collective exciters, which considerably simplified the flight tests. An amplitude of 100% equalled $\pm 5^{\circ}$ of flaperon motion.

An electronic controller automatically swept the exciters from 1 to 10 Hz, using a logarithmically increasing sweep rate of approximately ten cycles per octave. This was faster than the rate recommended in reference 9, but still slow enough to reveal each mode. Three such sweeps in succession were performed at each test condition.

The flight conditions tested are discussed later in this report under Flight Test Results.

ANALYSIS METHODS

The overall concept of the modal identification methods used in this study is to first estimate the frequency response H(f) between the aircraft excitation and structural response, and then to determine the modal damping and frequency by second-order model fitting. Two approaches for obtaining the frequency-response estimates are available; the choice of method depends upon the relative levels of measurable excitation (flaperon input) and unmeasurable excitation (turbulence), as discussed in reference 10.

Figure 5 illustrates the excitation of the aircraft by measurable and unmeasurable inputs x(t) and m(t); the measured response y(t) is corrupted by measurement noise n(t). If the measurable and unmeasurable inputs and measurement noise are fully uncorrelated, then the <u>unbiased</u> (true average value) frequency response H(f) may be estimated from the cross- and auto-spectral functions $G_{xy}(f)$ and $G_{xx}(f)$ as

$$H = \frac{G_{xy}}{G_{xx}} \tag{1}$$

(dependence on frequency f is implied). The random scatter in the estimate is a function of the noise-to-signal ratios G_{mm}/G_{xx} and G_{nn}/G_{vv} (see ref. 10). Thus, when sizeable, persistent, and measurable excitation of a mode is achieved (i.e., small noise-to-signal ratios), the calculation method of equation (1) is preferable; this (direct) approach is referred to herein as the "transfer-function method."

When adequate direct excitation cannot be achieved, as in the case of chord-mode response to flaperon inputs, an alternate method is used (fig. 6). In this (indirect) approach, referred to herein as the "cross-spectral method," structural responses from the right and left wings (y_R and y_L) are cross-correlated with each other, rather than with the flaperon inputs. The resulting cross-spectrum is

$$G_{y_R y_L} = H_R^* H_L [G_{xx} + G_{mm}]$$
 (2)

For symmetric modes, ${\rm H}_{\rm R}$ = ${\rm H}_{\rm L},$ and for antisymmetric modes, ${\rm H}_{\rm R}$ = ${\rm -H}_{\rm L}.$ In the symmetric case,

$$G_{y_R y_L} = |H|^2 [G_{xx} + G_{mm}]$$
 (3)

which allows an estimate of the frequency-response magnitude |H| (to within a scale factor) when the total excitation $[G_{\rm XX}+G_{\rm mm}]$ is constant within the frequency range of the modal response (which is generally a satisfactory assumption). In the cross-spectral method, the phase information in H is lost; however, if the noise-to-signal ratio $G_{\rm mm}/G_{\rm XX}$ is large, a much lower scatter in the magnitude estimate is achieved than would be possible with the transfer-function method. The damping

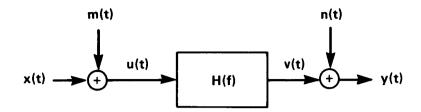


Figure 5.- Signals and noise affecting transfer-function calculations.

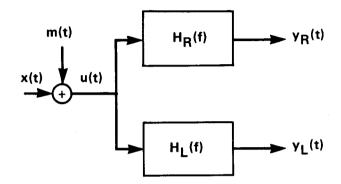


Figure 6.- Signals and noise affecting cross-spectrum calculations.

ratio and frequency are then determined from the magnitude response estimate (demonstrated in ref. 5).

Figure 7 schematically shows the procedures used to conduct the analyses discussed above. Each large block corresponds to a separate computer program. After each flight, the data are loaded into the Tilt-Rotor Engineering Database System (TRENDS) for ease of subsequent access. Next, the Frequency Response Identification (FRESPID) program generates the spectral functions from the time histories in TRENDS. Finally, the modal parameters are determined by using the curve-fitting program, NAVFIT. All computations are performed off-line (postflight).

TRENDS was developed by M. J. Bondi of NASA Ames Research Center and W. S. Bjorkman of Analytical Mechanics Associates, Inc. FRESPID was written by M. B. Tischler and J. G. M. Leung of Ames Research Center, and NAVFIT was originally developed by J. Hodgkinson and J. Buckley of McDonnell Douglas Aircraft (ref. 11). For a detailed discussion of FRESPID and NAVFIT, see reference 12.

The following sections briefly summarize the programs and discuss their application to XV-15 aeroelastics data.

Fourier-Transform Computations

The first step is to Fourier-transform the exciter and structural response data using FRESPID. For the transfer-function method, corresponding left and right transducers (strain gages or LVDTs) are summed or differenced, depending on the mode, to form composite inputs and outputs. If the two transducers are properly chosen, then the structural signals will be highly correlated and in phase for

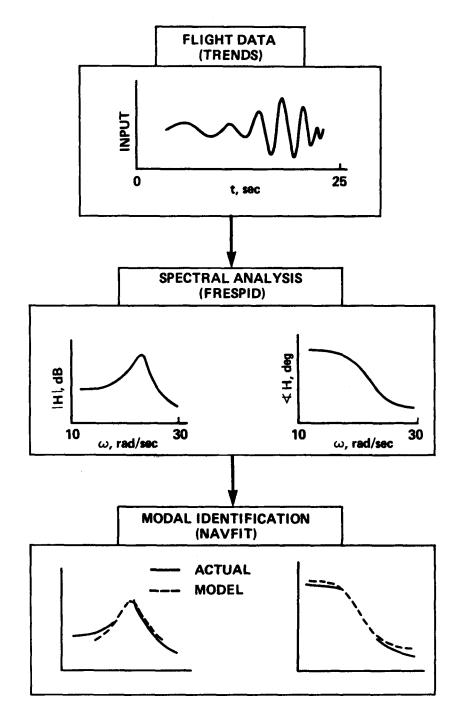


Figure 7.- Sequence of data processing for frequency-domain modal identification.

symmetric modes, and highly correlated but out of phase for antisymmetric modes. Noise will not be correlated; thus corruption of the spectral data is minimized.

FRESPID transforms the time-history data to the frequency domain by using a chirp z-transform. The chirp z-transform algorithm is an improvement on the

conventional Cooley-Tukey Fast Fourier Transform (FFT), in that it allows arbitrary resolution over a specified frequency range (ref. 13).

Figure 8 shows the computational procedures which are used by FRESPID to generate the auto- and cross-spectral functions from the input and output time-histories. The dc components and linear drifts are first removed to prevent oscillation in the spectral calculations. Multiple runs are concatenated to form composite time-histories. The (concatenated) time-histories are digitally filtered and partitioned into several overlapping sections. Each section is scaled with a cosine weighting function ("Hanning window") to prevent side lobes and leakage (ref. 14). The spectral content of each section is analyzed using the chirp z-transform. The total spectrum is finally determined by averaging the spectra of all of the sections.

Spectral Functions

Once the Fourier coefficients have been computed by the chirp z-transform, the auto- and cross-spectral functions, $G_{xx}(f)$, $G_{yy}(f)$, and $G_{xy}(f)$, and the transfer function, H(f), are calculated by the formulas in reference 10. The coherence function γ^2 is also computed; it is useful for evaluating the identification accuracy (ref. 10). The magnitudes of the spectral functions are converted to power decibels (e.g., 10 $\log_{10} |G_{xx}(f)|$) for plotting. The transfer-function results are presented in standard Bode form, that is, a semilog plot of magnitude (dB = 20 $\log_{10} |H(f)|$) and phase (deg) versus frequency (rad/sec).

Figure 9 shows the output time-history of the left wing-bending strain gage produced by a symmetric flaperon sweep (the right signal is nearly identical). The autospectrum G_{vx} corresponding to three such sweeps is fairly flat, as shown in figure 10. The associated cross-spectrum $G_{y_R}y_L$ between the right and left signals is plotted in figure 11, which clearly shows the peak of the first bending mode.

Figure 12 shows the transfer-function frequency and phase for the symmetric beam response to flaperon input. The shape of the magnitude peak of the mode is similar, but not identical, to that of the cross-spectrum (fig. 11), because the autospectrum is not precisely constant (fig. 10). The phase plot clearly shows the 90° change in phase at the natural frequency.

For the transfer-function method, the coherence function γ_{xy}^2 may be interpreted as that fraction of the output (response) spectrum that is linearly related to the input (excitation) spectrum (ref. 14). If the system is perfectly linear and noise-free, the coherence will be unity. For the cross-spectral method, high coherence γ_{xy}^2 implies a common excitation source to the right and left wings, and low output sensor noise. For both methods, the coherence is a good measure of the quality of the data prior to application of the modal curve fit.

Figure 13 illustrates the coherence functions corresponding to the cross-spectrum and transfer function shown in figures 11 and 12. The coherence associated with the cross spectrum is very nearly unity throughout the frequency range of the

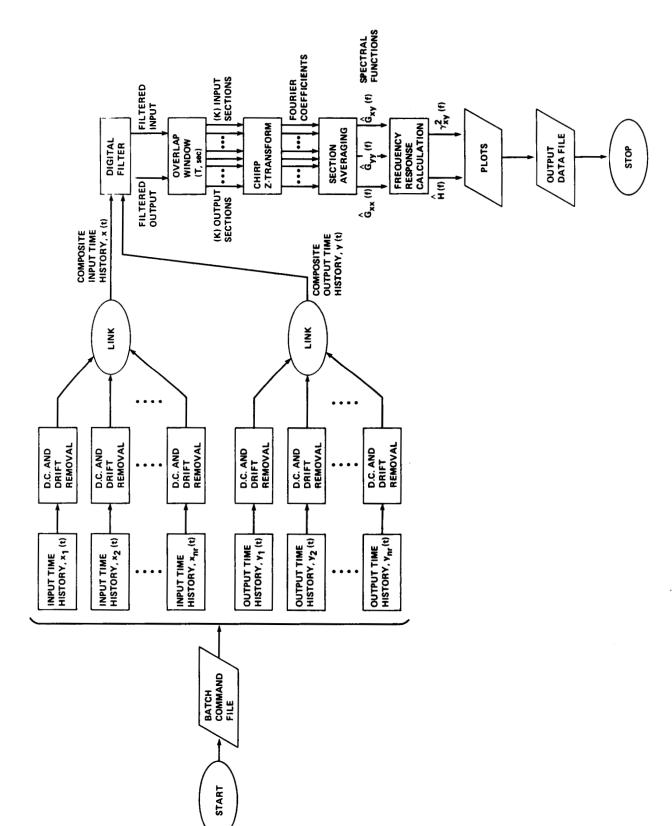


Figure 8.- Flowchart of computations performed by the FRESPID program.

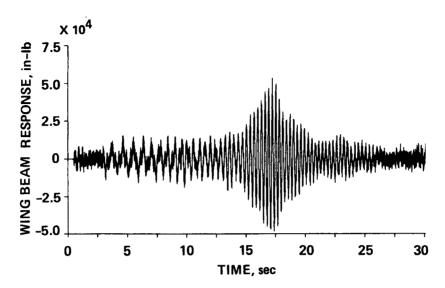


Figure 9.- Strain-gage output for one symmetric flaperon sweep (right-wing beam bending).

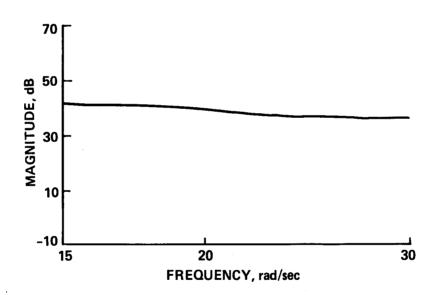


Figure 10.- Autospectrum G_{XX} of the summed flaperon inputs for three symmetric sweeps.

mode, indicating a highly linear response with little noise. However, the values for the transfer-function data are not as good as those for the cross-spectral data, which is a typical result. Worse coherence was often seen in other modes, especially the chord mode, where the response was excited less directly.

Frequency and Damping Calculations

Once cross-spectra and transfer functions have been calculated by the Fourier-transform program FRESPID, modal frequencies and damping are determined by curve-fitting the spectral data. An integration method was also tried for comparison.

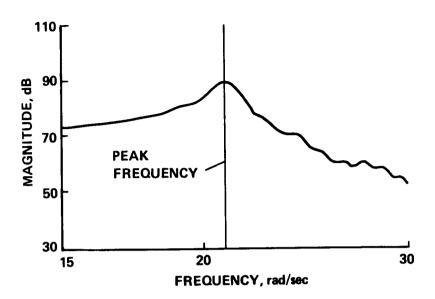


Figure 11.- Cross-spectrum $G_{y_Ry_L}$ of right and left wing-beam strain-gage outputs for three symmetric flaperon sweeps.

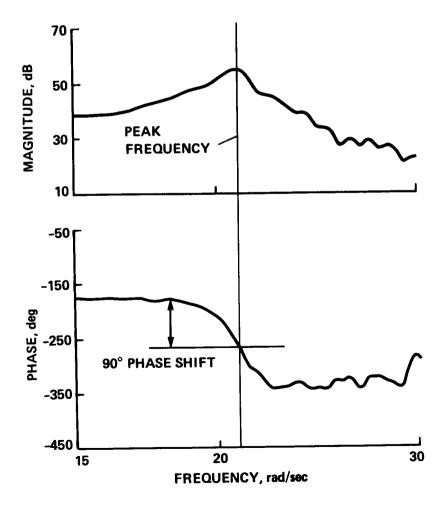


Figure 12.- Transfer function H (magnitude and phase) for the symmetric beam mode.

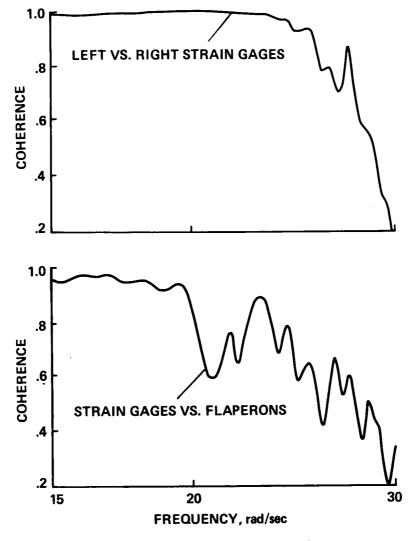


Figure 13.- Coherence functions $\gamma_{y_R}^2 y_L$ and γ_{xy}^2 associated with the dross-spectrum (fig. 11) and the transfer function (fig. 12).

The individual modes are well approximated by second-order quadratic models, which are functions of modal damping ζ and frequency ω_n . Only such models were used in the present study.

The curve-fitting program NAVFIT is a general multimode, high-order analysis using both magnitude and phase data. Either cross-spectral or transfer-function data can be analyzed with NAVFIT as desired. The user specifies a frequency range, the order of the model to be fitted, and (optionally) initial estimates of ζ and ω_n . Phase shifts caused by either unmodeled higher modes or 1/rev can be fitted with a time delay. An iterative algorithm is used to refine the model by systematically varying its coefficients. As discussed earlier, modal fitting is based on the magnitude and phase for the transfer-function method, but on magnitude alone for the cross-spectral method.

The model is fitted by minimizing a cost function based on the squares of both magnitude and phase errors. Frequency weighting is used to emphasize the data that are closest to the peak structural response. The relative weights of the magnitude and phase errors are chosen to yield results equivalent to equal weighting of the real and imaginary parts of the complex transfer function (ref. 12). For cross-spectral data, the phase weight is set to zero.

In addition to NAVFIT, a method was tried which computed the damping ratio from spectral data by integrating over the spectral peak. The integration program implemented the second-order system analysis discussed in reference 15 for cross-spectra and transfer functions. The algorithm was extended to allow two parameters, $\omega_{\rm n}$ and the peak spectral magnitude, to be iteratively varied to obtain a best fit to the data.

There were thus four possible calculations of frequency and damping: integration and curve fits, each applied with the cross-spectral and transfer-function methods. The final choice was based upon the scatter in estimating ζ and ω_n at the baseline point (defined below under Flight-Test Results). With one exception, the antisymmetric beam mode, NAVFIT gave the lowest scatter; even in the exceptional case, the NAVFIT results showed no statistically significant differences from the integration method.

The transfer-function method gave better results than the cross-spectral method, except for the chord modes. The flaperons cannot directly excite the wing in the chordwise direction as they can in bending and torsion. The resulting ratio of noise (atmospheric turbulence excitation) to measurable input is high, which explains why the cross-spectral method gives better results for chord modes. Various refinements to the transfer-function method are being explored in hopes of improving its noise rejection for the chord modes.

An example of the use of NAVFIT to determine frequency and damping for the symmetric beam mode is given in figure 14 (compare with fig. 12). Note that magnitude and phase are both fitted with a second-order response.

FLIGHT-TEST RESULTS

Figure 15 shows the portions of the XV-15 flight envelope covered during the aeroelastics flight tests. Because the aircraft had already been cleared to fly the envelope shown, the frequency sweeps were concentrated within a fairly narrow region so as to more rigorously verify the frequency-domain technique. The most comprehensive data were taken at 10,000-ft density altitude at 86% rotor speed (8.6 Hz). The speed range was 150 KIAS (Knots Indicated Airspeed), the normal speed for conversion to airplane mode, up to 220 KIAS, the torque-limited maximum speed for level flight. Only these data are reported here. (Limited data were also taken at 5,000 ft and 15,000 ft at 86% rotor speed, and at 10,000 ft at 98% rotor speed. Maximum-power climbs and power-off descents were performed at 150 KIAS at 86% rotor speed, with data taken as the aircraft passed 10,000 ft.)

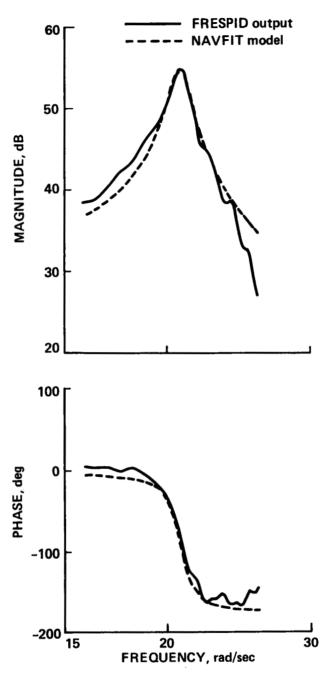


Figure 14.- Curve fits to transfer-function magnitude and phase for the symmetric beam mode.

Limitations on flight-test time did not permit replications of all test points. Therefore, a baseline point of 150 KIAS at 10,000 ft was chosen as an easily repeatable flight condition to explicitly test for scatter in the frequency and damping estimates.

Earlier flight tests (ref. 6) showed interaction of the Stability Control Augmentation System (SCAS) with modal responses, which was eliminated by

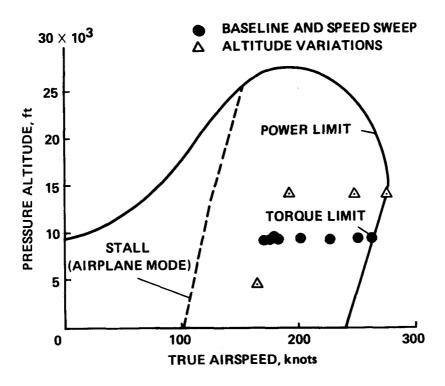


Figure 15.- XV-15 flight envelope, showing frequency-sweep test points.

modification of the SCAS. To ensure that there were no other interactions, each part of the automatic flight control system—the SCAS, the Attitude Retention System (ARS), and the Force Feel System (FFS)—was individually turned off during three series of sweeps at the baseline point. In a comparison of the results with the other baseline estimates, no statistically significant differences were noted. These data were subsequently included in the baseline data.

Summary Plots and Statistics

Figure 16 summarizes the frequency and damping results for all six modes. Open symbols are estimates made with exponential-decay methods (from ref. 16); closed symbols are frequency-domain estimates. (Reference 16 does not always give an explicit frequency estimate corresponding to each damping estimate.) The frequency-domain method yields low scatter at the baseline point and good consistency between airspeeds, hence higher reliability than the exponential-decay method. (Individual modes are discussed in detail below.)

A detailed assessment of all available predictive methods is beyond the scope of this paper. To illustrate discrepancies between predictions and flight data, representative results of two different analyses are shown in figure 16. The DYN4 predictions of frequency and damping were made by Bell Helicopters (ref. 16); the CAMRAD predictions were made by NASA. The CAMRAD model used here is that of reference 3, but it was modified to have the correct precone, rotor speed, and flight conditions. Neither analysis exactly models the aircraft. However, the trends in

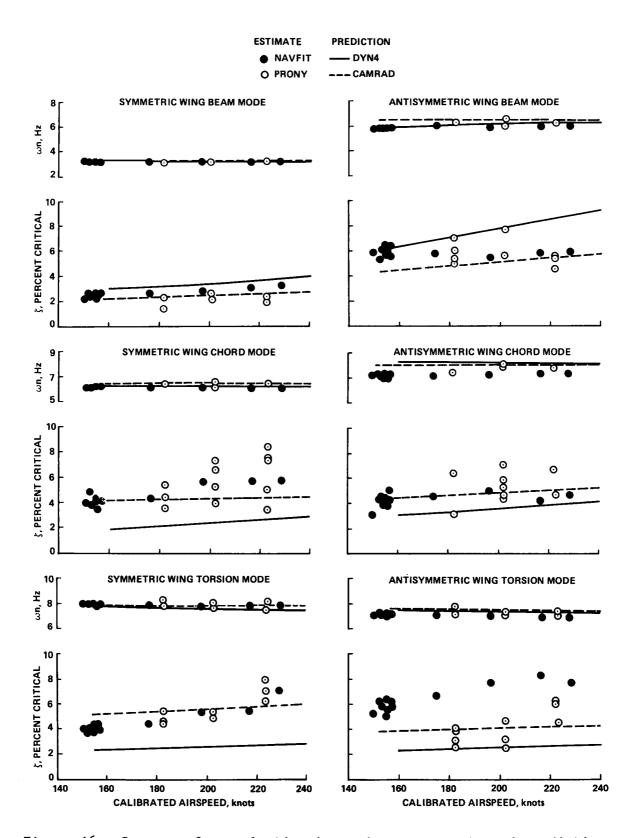


Figure 16.- Summary of aeroelastic wing-mode measurements and predictions.

the plotted results illustrate the differences between the two methods. Further-more, the DYN4 program does not represent the latest analytical techniques used by Bell Helicopters, but it is the method used during the design and initial flight test of the XV-15.

Numerical results of the frequency-domain method are summarized in Table 1 for the 150-KIAS baseline point. Listed for each mode are the averages of damping $(\bar{\zeta})$ and frequency $(\bar{\omega}_n)$ and their respective standard deviations (σ_ζ) and (σ_ω) , which are measures of the scatter among the estimates. The standard deviations of the damping range from 6% to 15% of the average values, while the standard deviations of the frequency are all less than 1%.

TABLE 1.- XV-15 WING MODES AT THE BASELINE FLIGHT CONDITION (STATISTICS ARE BASED ON 8 DATA POINTS)

Mode	ζ	σζ	σζ	ωn	σω	σ_{ω}
	% critical damping	<pre>\$ critical damping</pre>	% error	Hz	Hz	% error
Symmetric Beam	2.52	0.218	8.6	3.31	0.0107	0.32
Antisymmetric Beam	5.97	0.401	6.7	5.89	0.0427	0.73
Symmetric Chord	4.17	0.398	9.5	6.28	0.0198	0.32
Antisymmetric Chord	4.06	0.609	15.0	7.21	0.0440	0.61
Symmetric Torsion	4.08	0.245	6.0	8.07	0.0385	0.48
Antisymmetric Torsion	5.77	0.464	8.0	7.24	0.0349	0.48

In a very few cases--notably, antisymmetric torsion--a significant fraction of the scatter can be explained by weight changes caused by fuel burnoff. It is not practical to collect all flight data at exactly the same fuel state. Therefore, the values given in table 1 represent realistically achievable performance of the frequency-domain method.

There is no comprehensive set of exponential-decay data corresponding to the frequency-sweep baseline data. In order to make statistical comparisons between the results of the two flight-test methods, the frequency and damping results were curve-fitted against airspeed. This procedure allows consistent comparisons to be made using all of the data. Linear fits were used, partly because all predictions show nearly constant slopes within the airspeed range and partly because the standard error of each fit will be a conservative measure of scatter if the true variations are in fact nonlinear.

The standard errors, slopes, and intercepts of the damping curve fits are listed in table 2 for each mode. The standard error is a measure of the scatter about a fitted curve (analogous to the standard deviation about a point). The frequency-domain method generally has a standard error of one-half or less than the exponential-decay method.

TABLE 2.- CURVE-FIT STATISTICS OF DAMPING VERSUS AIRSPEED

		Frequency-Domain Method			Exponential-Decay Method*			
	No. of Points	Standard error, % critical	Intercept, % critical damping	Slope %/knot	No. of points	Standard error, % critical	Intercept, % critical damping	Slope %/k~ot
Symmetric Beam	12	.181	1.06	.00947	6	0.422	0.865	.0069
Antisymmetric Beam	12	.352	6.06	000776	9	1.02	8.82	015
Symmetric Chord	12	.405	0.442	.0405	12	1.64	-4.61	.050
Antisymmetric Chord	12	.565	2.84	.00823	9	1.64	0.883	.022
Symmetric Torsion	12	.332	-1.39	.0353	8	0.491	-5.36	.056
Antisymmetric Torsion	12	.446	0.557	.0341	10	1.01	-5.75	.049

*Based on data from reference 15

The key requirement is to detect any change in the trend of damping with airspeed as a stability boundary is approached. Accurate estimates of average damping are not as important as accurate estimates of slope, against which deviations can be detected. Most modes show increasing damping with airspeed, but will change to a negative slope at a sufficiently high airspeed (see refs. 3 and 6). For some modes, the negative slope is predicted to become very steep at the critical airspeed for flutter.

Figure 16 and tables 1 and 2 together show that the frequency-domain method is very reliable and completely adequate to detect a significant decrease in damping should a dangerous part of the flight envelope be entered. Even without further refinement of the flight-test or analytical techniques, the new approach is seen to be preferable to the exponential-decay method.

Care should be taken when comparing the exponential-decay results with the new frequency-sweep results. No complete set of the former type of data for an aircraft configuration exists that exactly matches that for the new data. The exponential-decay results shown in figure 16 are for 1.5° precone titanium hubs; the frequency-domain results are for 1.5° steel hubs. Both configurations are for the aircraft at NASA Ames with metal blades. These are the most closely-matched aeroelastics data sets available. Other reports on XV-15 aeroelastics (e.g., refs. 6, 7, and 16) sometimes include data for the aircraft operated by Bell Helicopters, which is not identical to the aircraft at NASA, or for 2.5° precone hubs. These additional configurations are thought to have slightly different aeroelastic behavior.

Individual Modes

For the <u>symmetric beam</u> mode, the new frequency-domain estimates of natural frequency ω_n are similar to the old exponential-decay results, and closely follow the predictions. The new estimates of damping ζ show a very consistent increase with airspeed. The CAMRAD estimate appears to match the new estimates of ω_n slightly better than DYN4, but the match for ζ is slightly worse. However, the

differences between the slopes of DYN4 and CAMRAD are too small to allow a reliable choice of one method over the other.

For the <u>antisymmetric beam</u> mode, the new estimates of ω_n show less scatter than the old estimates and are slightly lower in frequency than the predictions, especially for CAMRAD. The new estimates of ζ are much better than the old, but neither set of estimates matches the slope of either predictive method. However, CAMRAD does better approximate the measured slope of ζ .

For the <u>symmetric chord</u> mode, the new estimates of ω_n have slightly less scatter than the old, and are slightly lower in frequency than the predictions. The new estimates of ζ are much improved, but there is still noticeable scatter. There is too much scatter in even the new estimates to reliably favor either predictive method. It is not known whether the leveling off in the new estimates of ζ is an artifact of the scatter or an accurate measurement of a nonlinear variation with airspeed; more flight data are required to draw a clear conclusion.

For the <u>antisymmetric chord</u> mode, the new estimates of ω_n are slightly better, but some scatter remains. The estimates are about 0.5 Hz lower than the predictions and show an increase with airspeed not matched by the predictions. The new estimates of ζ are better, but scatter is noticeable. The slope in the estimates of ζ is not clear enough to reliably reveal any errors in the predictions.

For the <u>symmetric torsion</u> mode, the new estimates of ω_n are very good. They do not show the predicted decrease in ω_n as airspeed is increased from the baseline condition (this is more noticeable for DYN4 than for CAMRAD). The new estimates of ς are slightly better than the old. Both sets of estimates show increasing slope with airspeed, but it has not been proven that this indicates a real aeroelastic phenomenom. Neither DYN4 nor CAMRAD matches the slope of the estimates.

For the <u>antisymmetric torsion</u> mode, the new estimates of ω_n are definitely better. They are lower than the predicted values, but approximately match the slope predicted by CAMRAD. The new estimates of ζ are better than the old, and have almost twice the magnitude. Both predictions of ζ show lower slopes than the estimates. The new estimates show a decrease in slope at the last, highest-speed point, whereas the old estimates appear to show an increase in slope. It cannot be determined whether either case reflects a true change in slope or whether both are merely caused by scatter.

Overall, there is no obvious, consistent pattern of correspondence between either predictive method and the modal estimates based on flight data. The occasional large inconsistencies, known modeling errors (thought to be small), and newly available, improved analyses have prompted a rework of the aeroelastic predictions, which is now in progress. In any event, the excellent results of the frequency-domain analyses should allow reliable evaluations of new predictions.

CONCLUSIONS

Frequency-domain estimates of modal frequencies and damping were consistent with airspeed and were highly repeatable at a reference flight condition (within less than 1% for ω_n and 15% for ζ). These results demonstrate that the frequency-sweep method is a reliable and efficient way of determining XV-15 aeroelastic behavior from flight data. The ability to use either cross-spectral or transfer-function data provides great analytical flexibility, permitting good estimations of all modes. Such analytical methods can now be routinely used to support planned flight tests of the new XV-15 composite blades (ATBs) or any other new configuration.

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